

# Comparison of hydrogen sulfide gas and sulfur powder for synthesis of molybdenum disulfide nanosheets



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## ARTICLE INFO

### Article history:

Received 29 July 2015

Received in revised form

12 February 2016

Accepted 31 March 2016

Available online 1 April 2016

### Keywords:

Hydrogen sulfide

Sulfur powder

Transition metal dichalcogenides (TMDC)

2D nanosheet

Selectable area growth

## ABSTRACT

Layered transition metal dichalcogenides (TMDCs) have the potential to be used as an alternative to graphene in nano applications, because of their unique opto-electric properties. However, the large-scale synthesis of TMDCs has not been thoroughly investigated with different sulfurization agents, nor have the effects of defects and vacancies on such synthesis been determined. This study therefore looks at the synthesis of MoS<sub>2</sub> nanosheets from a thin Mo film to compare the effects of using H<sub>2</sub>S gas or sulfur powder as the sulfurization agent. In either case, a four-layered nanosheet is obtained; however, the good stoichiometry of MoS<sub>2</sub>, smallest defects, and fewest vacancies are obtained using H<sub>2</sub>S.

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## 1. Introduction

Layered transition metal dichalcogenides (TMDCs) have recently drawn significant attention for use as next generation nano-electronics [1], thin film transistors [2], logic circuits [3], sensors [4], catalysts [5,6], and optoelectronic devices [7,8]. The physical, electrical, and optical properties of TMDCs have so far mostly been studied through physical exfoliation of nanosheets and their transfer onto suitable substrates [9–11]. However, nanosheets with a large surface area can only be achieved through synthesis directly onto a substrate. Recently, the layered MoS<sub>2</sub> was most famous in TMDCs because of the exotic properties of optical and electrical. The early stage method of synthesizing MoS<sub>2</sub> is the sulfurization of metal and metal-oxide films using a sulfurization agent [12–16], and good quality MoS<sub>2</sub> and WS<sub>2</sub> nanosheets have been achieved through chemical vapor deposition (CVD) [17–20].

There is a potential problem in the use of sulfur powder as a sulfurization agent, namely the fact that its vaporization requires temperatures above 200 °C, and that its vapor has a S<sub>8</sub> crown ring

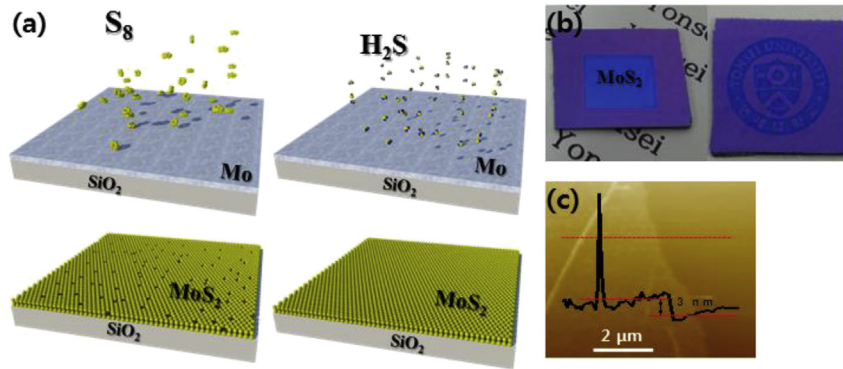
structure that leads to inhomogeneous sulfurization [21]. In contrast, H<sub>2</sub>S gas has a simple molecular structure with a low decomposition energy, ensuring a reliable and homogeneous reaction [22,23]. The synthesis of MoS<sub>2</sub> nanostructures has been previously achieved using H<sub>2</sub>S gas [24, 25], but was found to result in S vacancies forming near nanocluster edges, grain boundaries, and defect sites [26, 27, 28] that greatly affected the optical and electrical properties [27, 28]. However, as there has been no comparative study of different sulfurization agents, this study compares the synthesis of MoS<sub>2</sub> using H<sub>2</sub>S gas with that using sulfur powder at temperatures of 1000 and 750 °C. The intrinsic properties of the MoS<sub>2</sub> nanosheets obtained are also investigated through pump probe measurements, though the ease with which MoS<sub>2</sub> readily changes its electrical properties prevented these from being included.

## 2. Experimental section

The single-crystal SiO<sub>2</sub> and MgO substrates were prepared by ultrasonic cleaning in acetone, IPA, and DI water for 10 min each. For metal sulfurization, 5 or 10 nm Mo metal films were deposited by sputtering for 5 or 10 s, respectively. Fig. 1 (a) shows the

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**Fig. 1.** (a) Schematic illustration of the sulfurization process. (b) Optical images of MoS<sub>2</sub> layers grown on a SiO<sub>2</sub> substrate. (c) AFM image of a four-layer MoS<sub>2</sub> nanosheet ~4 nm in thickness on a SiO<sub>2</sub>/Si substrate.

differences in the sulfurization process depending on whether initial sulfurization in a tube type furnace at 500 °C was carried out using 5 sccm of H<sub>2</sub>S gas or 3 g of sulfur powder (introduced at 120 °C) under a 50 sccm flow of Ar. Following this initial sulfurization, the temperature was increased to 750 or 1000 °C and held for 10 min to enhance the quality of the MoS<sub>2</sub> nanosheet. The sample preparation conditions are summarized in Fig. 2.

To measure the thickness of the MoS<sub>2</sub> nanosheets produced, each was measured by atomic force microscopy (AFM). Differences between the two different sulfurization methods were assessed by characterizing the different nanosheets using Raman spectroscopy, X-ray photoemission spectroscopy (XPS), tunneling electron microscopy (TEM), and pump probe measurement. The XPS measurements were performed at the 09A beam line of the National Synchrotron Radiation Research Center (NSRRC) in Hsinchu, Taiwan, with all data being obtained at a base pressure of less than  $1 \times 10^{-9}$  mbar.

The electrical characteristics and physical properties of the MoS<sub>2</sub> nanosheet were analyzed using ultrafast pump-probe methods so as to prevent any contamination or damage. These measurements were carried out with a visible probe range based on a 250 KHz Ti:sapphire laser system (Coherent RegA 9050) producing 50 fs pulses at 800 nm. A white-light continuum was created by focusing these 800 nm pulses into a sapphire disk, with 400 nm pulses being generated by a second harmonic.

### 3. Results and discussion

As shown in Fig. 1(b), large-area MoS<sub>2</sub> nanosheet with a uniform color was selectively grown using H<sub>2</sub>S gas. After transferring it onto a clean SiO<sub>2</sub> substrate, an AFM line profile (Fig. 1(c)) was used to determine its thickness as 3 nm, which equates to four layers of MoS<sub>2</sub> nanosheet.

Fig. 2(a) shows the Raman spectra of MoS<sub>2</sub> nanosheets, prepared under various conditions (Fig. 2), from which the width of each peak can be used to determine the crystallinity of the sample [29, 30]. That is, the  $24 \text{ cm}^{-1}$  gap between the E<sub>12g</sub> and A<sub>1g</sub> peaks corresponds to four layers of the MoS<sub>2</sub> nanosheet [30], and therefore confirms the AFM-measured thickness. Note also that the Si peak is normalized to allow such comparison between the intensities and widths of the E<sub>12g</sub> and A<sub>1g</sub> peaks. Fig. 2(b) shows the peak distance as a function of the Mo film thickness, suggesting a strong relation between the thickness of the MoS<sub>2</sub> nanosheets and the Mo films. There is, however, a poor conformity to the Raman spectrum when the number of layers exceeds four, with the larger intensity ratio of sample H10 indicating that a good crystallinity can be obtained by annealing at 1000 °C with H<sub>2</sub>S gas. The full-width at half maximum

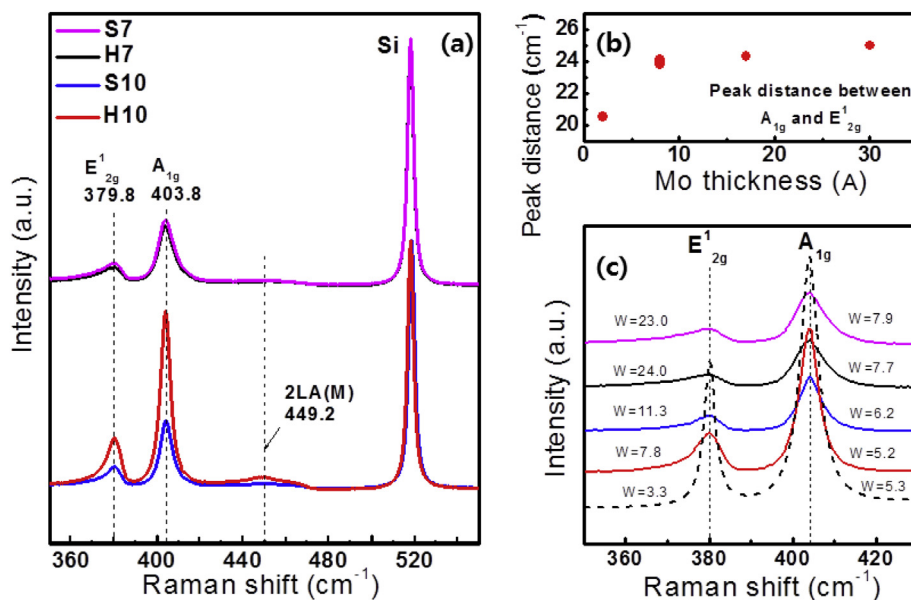
(FWHM) values of the Raman peak given in Fig. 2(c) reveal that the width of the E<sub>12g</sub> peak relative to that of the A<sub>1g</sub> peak is lower in the synthesized MoS<sub>2</sub> nanosheet than in mechanically exfoliated-MoS<sub>2</sub> (the black dotted line), which can be attributed to planar vibration (E<sub>12g</sub>) being strongly related to the randomly distributed grains found in CVD-grown MoS<sub>2</sub> nanosheets [13]. The FWHM values and peak intensities are also much higher at 750 °C than at 1000 °C, especially when H<sub>2</sub>S gas is used for sulfurization.

The quality and crystallinity of the MoS<sub>2</sub> nanosheet were assessed by TEM, with Fig. 3(a,b) showing the nanosheet prior to sulfurization. The use of H<sub>2</sub>S (Fig. 3(d,f)) was found to produce an enhanced crystallinity when compared to using sulfur powder at the same annealing temperature (Fig. 3(c,e)). The diffraction patterns and FWHM in the insets of Fig. 3(c–f) show evidence of a six-fold-symmetry following sulfurization [13, 31], with the ring pattern of sample H10 indicating a polycrystalline structure.

The XPS results for the Mo 3d and S 2p core-levels of samples S7, H7, S10, and H10 are shown in Fig. 4, from which the atomic ratio (stoichiometry) of each sample was obtained based on the ratio of the Mo 3d peaks intensity to that of the S 2p peak. As shown in Fig. 2, samples prepared in H<sub>2</sub>S gas had a comparable atomic concentration to bulk MoS<sub>2</sub>, whereas the use of sulfur powder resulted in concentrations either under (750 °C) or over (1000 °C) stoichiometry. This non-stoichiometry is believed to be caused by vacancies and defects generated during synthesis [26, 27, 28]. With either sulfurization agent, the width of the Mo 3d and S 2p core-level peaks roughly doubled; however, with only sulfur powder, the Mo 3d core levels exhibited a peak suggesting oxidation to MoO<sub>3</sub> at a higher binding energy. This phenomenon most likely originates from the substrate, with the oxygen species produced believed to interrupt the crystal structure and degrade the electrical properties of the thin film. This suggests that the use of H<sub>2</sub>S gas is the more suitable option for the synthesis of MoS<sub>2</sub> nanosheets.

To further confirm the superiority of H<sub>2</sub>S gas over sulfur powder, ultrafast optical-pump-probe measurements of the carrier relaxation time were performed. For this, four different samples of MoS<sub>2</sub> nanosheet were prepared on MgO substrates under conditions of: H<sub>2</sub>S at 750 °C (Sample I), sulfur powder at 750 °C (Sample II), H<sub>2</sub>S gas at 1000 °C (Sample III), and sulfur powder at 1000 °C (Sample IV). A MgO (100) substrate was used as this is not excited by probing measurements with a large band gap of 7.8 eV, as transparent substrates need a band gap larger than the photon energy (3.01 eV) of the pump pulses used for measurement.

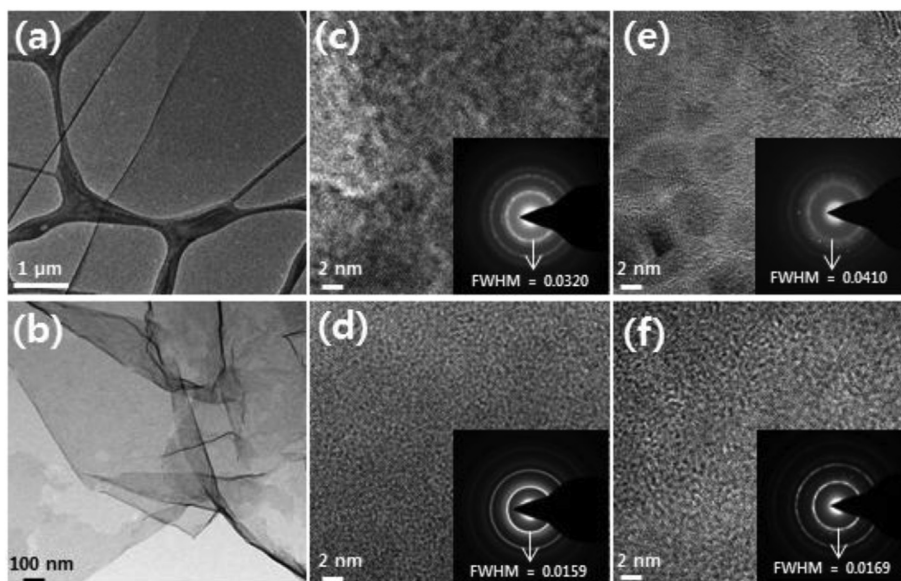
Prior to pump-probe measurement, the ground state absorption was measured to determine an appropriate pump-probe wavelength. As can be seen in Fig. 5 (a), the A and B excitons are



**Table 1.**

Sample name	temperature	Sulfur source	Sample name	temperature	Sulfur source
S7	750	S powder	H7	750	H <sub>2</sub> S gas
S10	1000	S powder	H10	1000	H <sub>2</sub> S gas

**Fig. 2.** Raman spectra of MoS<sub>2</sub> nanosheets produced using H<sub>2</sub>S and S at annealing temperatures of 750 and 1000 °C (a),(c). (b) Peak distance of each sputtered Mo film.

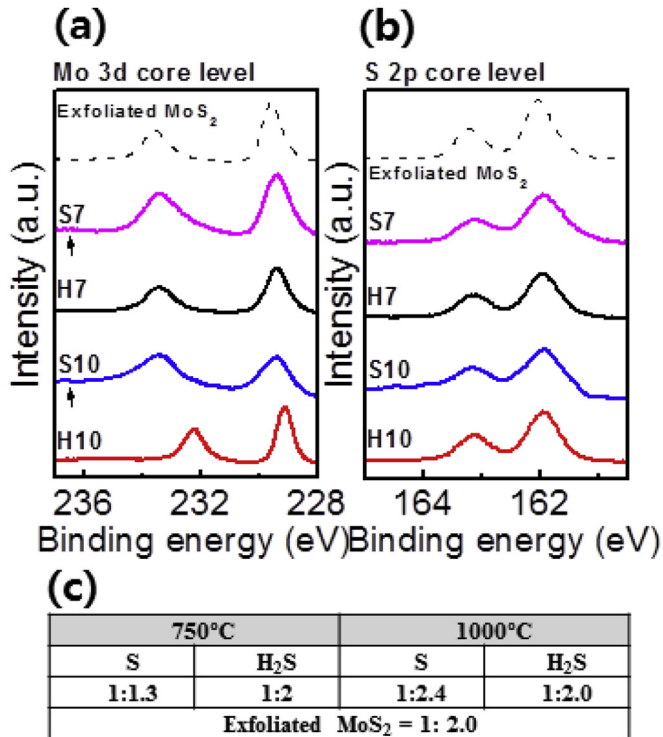


**Fig. 3.** Low-magnification TEM images of MoS<sub>2</sub> nanosheets (a,b) after transfer, and following sulfurization with (c) H<sub>2</sub>S at 750 °C, (d) sulfur powder at 750 °C, (e) H<sub>2</sub>S at 1000 °C, and (f) sulfur powder at 1000 °C. Insets show the respective FWHM (nm<sup>-1</sup>) and diffraction patterns.

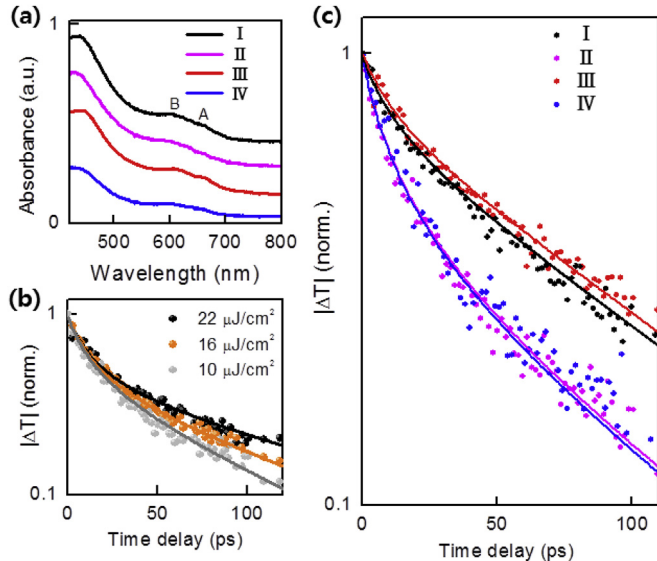
associated with direct transitions from a split valence band that results from spin-orbit splitting in the MoS<sub>2</sub> interlayer, with any interaction further modifying this split [32]. All samples exhibit resonance in the 600–700 nm range, which is well known as the direct excitonic transition position of the k point of the Brillouin zone [33]. The transition between peaks A and B [33, 34] is more resolved in samples I and III, with the appearance of these two peaks suggesting a low level of disorder in the sample [34]. This matches well with the earlier assessment that H<sub>2</sub>S gas produces

fewer vacancies, and from the ground–state absorption, a probe wavelength of 660 nm used in previous ultrafast pump-probe analyses of MoS<sub>2</sub> thin films [32, 35] was deemed appropriate.

A vacancy is a type of point defect in a crystal, with the scattering of carriers by impurities and lattice defects contributing to their relaxation [36]. Thus, in the pump-probe measurements, the changes in the decay dynamics of the pump-fluence were first observed for the main mechanism of carrier relaxation. All samples were found to show a tendency toward decreasing carrier



**Fig. 4.** XPS spectra of MoS<sub>2</sub> nanosheets sulfurized with H<sub>2</sub>S or S at 750 or 1000 °C. Peaks for (a) Mo 3d and S 2s, and (b) S 2p are scaled for clarity. (c) Chemical composition of each MoS<sub>2</sub> nanosheet.



**Fig. 5.** Ultrafast carrier relaxation dynamics for four different MoS<sub>2</sub> samples. (a) Ground state absorption spectra. (b) Normalized differential transmission signal (dots) of sample II using pump and probe beam wavelengths of 400 and 660 nm, respectively. (c) Normalized differential transmission signals (dots) of samples I, II, III, and IV with the same pump and probe beam wavelengths as (b). A fixed pump fluence of 16 J/cm<sup>2</sup> was used, and all solid lines are numerical fits obtained using rate equations that describe the vacancy-mediated pump fluence.

relaxation time with increasing pump-fluence; indeed, Sample (Fig. 5(b)) clearly exhibited a slow relaxation behavior with increasing pump-fluence. In a previous report, this slow relaxation was considered to originate from a limited defect state that is

rapidly saturated with carriers under the effects of a high fluence [37, 38]. We can therefore assume that the rapid decay and saturation from defects is the major factor in carrier relaxation. In Fig. 5 (c), numerical data for the samples were used to produce rate-equations for a fixed pump-fluence [37] based on:

$$\frac{dn(t)}{dt} = -\frac{n(t)}{\tau} + \frac{dT(t)}{dt}$$

$$n(0) = N, T(0) = T$$

where  $n(t)$  is the density of mobile charge carriers,  $\gamma$  (= 80 ps) is the carrier life-time,  $N$  is the initial density of free charges,  $T(t)$  is the density of unoccupied trap states, and  $T$  is the initial density of unoccupied trap states. These rate-equations demonstrate that the density of unoccupied trap states reduces over time due to a reduction in the density of mobile charge carriers, with the lifetime of any given defect state  $n(t)$  determined by the initial density of unoccupied trap states  $T$ . The initial density of free charge  $N$  of  $2 \times 10^{11}$  cm<sup>-2</sup> generated per layer with a pump-fluence of 16 J/cm<sup>2</sup> is obtained by using the method suggested by Shi et al. [32] From the fitted results in Fig. 5(c), it is clear that the relaxation dynamics of a carrier are not influenced by the annealing temperature; however, the carrier life-time is longer when using H<sub>2</sub>S instead of sulfur. The initial density of unoccupied trap states in Sample A, B, C, and D was 28, 55, 25, and 56, respectively, indicating that the number of defects is effectively halved by using H<sub>2</sub>S. Thus, the change in defects and vacancies is substantially greater than the change in crystallinity due to temperature.

#### 4. Conclusion

This study has demonstrated that the sulfuration can provide effective control over the thickness and area of a MoS<sub>2</sub> nanosheet. The electrical properties of the nanosheet were found to be closely linked to the number of defects and vacancies, which in turn were affected by the sulfuration agent used to a far greater extent than is seen in CVD. This means that a better outcome in terms of ensuring a low sulfur vacancy and high crystallinity can be achieved by using H<sub>2</sub>S gas rather than sulfur powder. It is hoped that these results will further the development of 2D-TMDC nanosheets and lead to more advanced sulfuration processes.

#### Acknowledgements

This research was supported by Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Education, Science and Technology (No. 2015R1D1A1A01060064), the National Research Foundation of Korea (NRF) grant funded by the Korea government (MSIP) (No. NRF-2014R1A2A1A11052588), the Center for Integrated Smart Sensors funded by the Ministry of Science, ICT & Future Planning as Global Frontier Project (CISS-2011-0031848), Korea Evaluation Institute of Industrial Technology (KEIT) funded by the Ministry of Trade, Industry and Energy (MOTIE) (Project No. 10050296, Large scale (Over 8'') synthesis and evaluation technology of 2D chalcogenides for next generation electronic devices), and the Institute of BioMed-IT, Energy-IT and Smart-IT Technology (BEST), a Brain Korea 21 plus program, Yonsei University.

#### References

- [1] M. Xu, T. Liang, M. Shi, H. Chen, *Chem. Rev.* 113 (2013) 3766.
- [2] S. Kim, A. Konar, W.-S. Hwang, J.H. Lee, J. Lee, J. Yang, C. Jung, H. Kim, J.-B. Yoo, J.-Y. Choi, Y.W. Jin, S.Y. Lee, D. Jena, W. Choi, K. Kim, *Nat. Commun.* 3 (2012) 1011.

- [3] H. Wang, L. Yu, Y.-H. Lee, Y. Shi, A. Hsu, M.L. Chin, L.-J. Li, M. Dubey, J. Kong, T. Palacios, *Nano Lett.* 12 (2012) 4674.
- [4] F.K. Perkins, A.L. Friedman, E. Cobas, P.M. Campbell, G.G. Jernigan, B.T. Jonker, *Nano Lett.* 13 (2013) 668.
- [5] D. Voiry, H. Yamaguchi, J. Li, R. Silva, D.C.B. Alves, T. Fujita, M. Chen, T. Asefa, V.B. Shenoy, G. Eda, M. Chhowalla, *Nat. Mater.* 12 (2013) 850.
- [6] J. Kibsgaard, Z. Chen, B.N. Reinecke, T.F. Jaramillo, *Nat. Mater.* 11 (2012) 963.
- [7] H.S. Lee, S.-W. Min, Y.-G. Chang, M.K. Park, T. Nam, H. Kim, J.H. Kim, S. Ryu, S. Im, *Nano Lett.* 12 (2012) 3695.
- [8] N. Perea-López, A.L. Elías, A. Berkdemir, A. Castro-Beltran, H.R. Gutiérrez, S. Feng, R. Lv, T. Hayashi, F. López-Urías, S. Ghosh, B. Muchharla, S. Talapatra, H. Terrones, M. Terrones, *Adv. Funct. Mater.* 23 (2013) 5511.
- [9] J.S. Ross, S. Wu, H. Yu, N.J. Ghimire, A.M. Jones, G. Aivazian, J. Yan, D.G. Mandrus, D. Xiao, W. Yao, X. Xu, *Nat. Commun.* 4 (2013) 1474.
- [10] Y. Li, C.-Y. Xu, L. Zhen, *Appl. Phys. Lett.* 102 (2013) 143110.
- [11] K. Lee, H.-Y. Kim, M. Lotya, J.N. Coleman, G.-T. Kim, G.S. Duesberg, *Adv. Mater.* 23 (2011) 4178.
- [12] Y.-C. Lin, W. Zhang, J.K. Huang, K.K. Liu, Y.H. Lee, C.T. Liang, C.W. Chu, L.J. Li, *Nanoscale* 4 (2012) 6637.
- [13] Y. Zhan, Z. Liu, S. Najmaei, P.M. Ajayan, J. Lou, *Small* 8 (2012) 966.
- [14] H. Hadouda, J. Pouzeta, J.C. Bernedeá, A. Barreaub, *Mater. Chem. Phys.* 42 (1995) 291.
- [15] X. Wang, H. Feng, Y. Wu, L. Jiao, *J. Am. Chem. Soc.* 135 (2013) 5304.
- [16] J.-G. Song, J. Park, W. Lee, T. Choi, H. Jung, C.W. Lee, S.-H. Hwang, J.M. Myoung, J.-H. Jung, S.-H. Kim, C. Lansalot-Matras, H. Kim, *ACS Nano* 7 (2013) 11333.
- [17] Y.-H. Lee, L. Yu, H. Wang, W. Fang, X. Ling, Y. Shi, C.-T. Lin, J.-K. Huang, M.-T. Chang, C.-S. Chang, M. Dresselhaus, T. Palacios, L.-J. Li, J. Kong, *Nano Lett.* 13 (2013) 1852.
- [18] S. Balendhran, J.Z. Ou, M. Bhaskaran, S. Sriram, S. Ippolito, Z. Vasic, E. Kats, S. Bhargava, S. Zhuiykov, K. Kalantar-Zadeh, *Nanoscale* 4 (2012) 461.
- [19] S. Najmaei, Z. Liu, W. Zhou, X. Zou, G. Shi, S. Lei, B.I. Yakobson, J.-C. Idrobo, P.M. Ajayan, J. Lou, *Nat. Mater.* 12 (2013) 754.
- [20] A.M. van der Zande, P.Y. Huang, D.A. Chenet, T.C. Berkelbach, Y.M. You, G.-H. Lee, T.F. Heinz, D.R. Reichman, D.A. Muller, J.C. Hone, *Nat. Mater.* 12 (2013) 554.
- [21] D.R. Alfonso, *Surf. Sci.* 602 (2008) 2758.
- [22] S. Wang, C. An, J. Yuan, *Materials* 3 (2010) 401.
- [23] J.A. Rodriguez, S. Chaturvedi, M. Kuhn, J. Hrbek, *J. Phys. Chem. B* 102 (1998) 5511.
- [24] X.L. Li, J.P. Ge, Y.D. Li, *Chem. Eur. J.* 10 (2004) 6163.
- [25] Y. Feldman, E. Wasserman, D.J. Srolovitz, R. Tenne, *Science* 267 (1995) 222.
- [26] W. Zhou, X. Zou, S. Najmaei, Z. Liu, Y. Shi, J. Kong, J. Lou, P.M. Ajayan, B.I. Yakobson, J.-C. Idrobo, *Nano Lett.* 13 (2013) 2615.
- [27] S. Tongay, J. Suh, C. Ataca, W. Fan, A. Luce, J.S. Kang, J. Liu, C. Ko, R. Raghunathanan, J. Zhou, F. Ogletree, J. Li, J.C. Grossman, *J. Wu, Sci. Rep.* 3 (2013) 2657.
- [28] S. Helveg, J.V. Lauritsen, E. Laegsgaard, I.I. Stensgaard, J.K. Nørskov, B.S. Clausen, H. Topsøe, F. Besenbacher, *Phys. Rev. Lett.* 84 (2000) 951.
- [29] C. Lee, H. Yan, L.E. Brus, T.F. Heinz, J. Hone, S. Ryu, *ACS Nano* 4 (2010) 2695.
- [30] S.-L. Li, et al., *ACS Nano* 6 (2012) 7381.
- [31] Y.-H. Lee, H. Miyazaki, H. Song, H. Kuramochi, S. Nakaharai, K. Tsukagoshi, *Adv. Mater.* 24 (2012) 2320.
- [32] H. Shi, R. Yan, S. Bertolazzi, J. Brivio, B. Gao, A. Kis, D. Jena, H.G. Xing, L. Huang, *ACS Nano* 7 (2013) 1072.
- [33] K.F. Mak, C. Lee, J. Hone, J. Shan, T.F. Heinz, *Phys. Rev. Lett.* 105 (2010) 136805.
- [34] G. Eda, H. Yamaguchi, D. Voiry, T. Fujita, M. Chen, M. Chhowalla, *Nano Lett.* 11 (2011) 5111.
- [35] R. Wang, B.A. Ruzicka, N. Kumar, M.Z. Bellus, H.-Y. Chiu, H. Zhao, *Phys. Rev. B* 86 (2012) 045406.
- [36] K. Ishioka, M. Hase, M. Kitajima, K. Ushida, *Appl. Phys. Lett.* 78 (2001) 3965.
- [37] P. Parkinson, H.J. Joyce, Q. Gao, H.H. Tan, X. Zhang, J. Zou, C. Jagadish, L.M. Herz, M.B. Johnston, *Nano Lett.* 9 (2009) 3349.
- [38] K. Lui, F. Hegmann, *J. Appl. Phys.* 93 (2003) 9012.